

# Noise-assisted energy transport in electrical oscillator networks with off-diagonal dynamical disorder

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Noise is generally thought as detrimental for energy transport in coupled oscillator networks. However, it has been shown that for certain coherently evolving systems, the presence of noise can enhance, somehow unexpectedly, their transport efficiency; a phenomenon called environment-assisted quantum transport (ENAQT) or dephasing-assisted transport. Here, we report on the experimental observation of such effect in a network of coupled electrical oscillators. We demonstrate that by introducing stochastic fluctuations in one of the couplings of the network, a relative enhancement in the energy transport efficiency of  $22.5 \pm 3.6\%$  can be observed.

Transport phenomena are ubiquitous throughout different fields of research. Some of the most common examples of transport analysis are seen in the fields of physics, chem-

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istry and biology.<sup>1</sup> In recent years, energy transport assisted by noise<sup>2-4</sup> has attracted a great deal of attention, partly because of its potential role in the development of future artificial light-harvesting technologies.<sup>5-7</sup> This intriguing phenomenon has theoretically been shown to occur in several quantum<sup>8-14</sup> and classical<sup>15-18</sup> systems; however, efforts towards its experimental observation had not been presented until very recently. Viciani *et al.*<sup>19</sup> showed an enhancement in the energy transport of optical fiber cavity networks, where the effect of noise on the system was introduced by averaging the optical response of several network configurations with different cavity-frequency values. In a closely related experiment, Biggerstaff *et al.*<sup>20</sup> demonstrated an increase in the transport efficiency of a laser-written waveguide network, where decoherence effects were simulated by averaging the output signal of the waveguide array considering different illumination wavelengths. Using the same photonic platform, Caruso *et al.*<sup>21</sup> observed an enhanced transport efficiency when suppressing interference effects in the transport dynamics of a photonic network. In this experiment, noise was implemented by dynamically modulating the propagation constants of the waveguides, which is the natural way for producing decohering noise, as it has been experimentally demonstrated in the context of quantum random walks.<sup>22-24</sup>

In this work, we report on the observation of noise-assisted energy transport in a network of capacitively coupled  $RLC$  oscillators, where  $R$  stands for resistance,  $L$  for inductance, and  $C$  for capacitance. Although in previous studies of ENAQT noise has been modeled as fluctuations in the frequency of each oscillator, so-called diagonal fluctuations,<sup>2</sup> here we introduce noise in the system by means of stochastic fluctuations in one of the network's capacitive couplings, referred to as off-diagonal dynamical disorder.<sup>25</sup> Using this system, we show that fluctuations in the coupling can indeed influence the system so that the energy transferred to one of the oscillators is increased, demonstrating that off-diagonal dynamical disorder can effectively be used for enhancing the efficiency of energy transport systems.

## Results

We consider a network of three identical *RLC* oscillators (as shown in Fig. 1), whose dynamics are described by

$$\frac{dV_n}{dt} = -\frac{1}{C} \left[ i_n + \frac{V_n}{R} + \sum_{m \neq n}^3 C_{nm} \left( \frac{dV_n}{dt} - \frac{dV_m}{dt} \right) \right], \quad (1)$$

$$\frac{di_n}{dt} = \frac{V_n}{L}, \quad (2)$$

where  $V_n$  is the voltage in each oscillator,  $i_n$  is the current, and  $C_{nm}$  represents the capacitive couplings.

Noise is introduced in the system by inducing random fluctuations in one of the capacitive couplings, so that

$$C_{12}(t) = C_{12} [1 + \phi(t)], \quad (3)$$

with  $C_{12}$  being the average capacitance of the coupling and  $\phi(t)$  a Gaussian random variable with zero average, i.e.  $\langle \phi(t) \rangle = 0$ , where  $\langle \dots \rangle$  denotes stochastic averaging.

Previous studies of noise-assisted transport have shown that efficiency enhancement can be observed by measuring the energy that is irreversibly dissipated in one particular site of the network, the so-called sink or reaction center.<sup>3,8,14,17</sup> Here, we take the resistance in Oscillator-2 to be the sink and measure the relative energy that is dissipated through it by computing

$$W = \frac{1}{\mathcal{N}} \int_0^\infty \left\langle \frac{V_2^2(t)}{R} \right\rangle dt, \quad (4)$$

where  $\mathcal{N}$  is the total energy dissipated by all the oscillators. Notice that Eq. (4) is equivalent to the efficiency measure that was derived in a previous work on noise-assisted transport in classical oscillator systems.<sup>17</sup>

We have experimentally implemented the system described by Eqs. (1)–(3) using functional blocks synthesized with operational amplifiers and passive linear electrical components (see Methods and accompanying Supplementary Information). Our experimental setup was designed so that the frequencies of each oscillator were the same ( $\nu = 290.57$  Hz), as well as the couplings between them ( $C_{nm} = 40 \mu\text{F}$ ). Notwithstanding, based on our measurements, we found that the designed system was characterized by the following parameters:  $\nu = 283.57$  Hz,  $C_{12} = 39 \mu\text{F}$ ,  $C_{13} = 41.7 \mu\text{F}$ , and  $C_{23} = 39 \mu\text{F}$ . These

variations in the parameters of the system result from the tolerances (around 5%) of all the electronic components used in the implementation of the circuit.

As described in the methods section, noise was introduced in one of the network's capacitive couplings using a random signal provided by a function generator (Diligent Analog Discovery 410-244P-KIT). The electronic circuit was designed so the voltage of the noise signal,  $V_s$ , is directly map into Eq. (3), thus making the stochastic variable  $\phi(t)$  fluctuate within the interval  $(-V_s, +V_s)$ , with a 1 kHz frequency. Notice that the frequency of the noise signal is higher than the oscillators' natural frequency, which guarantees a true dynamical variation of the capacitive coupling. Figure 2 shows some examples of the histograms of noise signals extracted directly from the function generator.

Using the configuration described above, we measured the energy dissipated through the resistor in Oscillator-2 [Eq. (4)] as a function of the noise voltage introduced in the capacitive coupling. We can see from Fig. 3 that transport efficiency in Oscillator-2 is enhanced as the noise voltage increases, a sign of noise-assisted energy transport.<sup>17</sup> We obtained an enhancement of  $22.50 \pm 3.59\%$  for a maximum noise voltage of 850 mV. It is important to remark that we cannot go beyond this value using the present configuration because, when introducing noise signals close to 1 volt (or more), the system becomes unstable due to the presence of negative capacitances in the coupling [see Eq. (3)]. However, as obtained in our numerical simulations, a higher efficiency may be reached by incorporating random fluctuations in the remaining couplings.

To verify that the observed enhancement was a consequence of energy rearrangement due to random fluctuations in the coupling, and not because external energy was introduced in the electronic circuit, we measured the transport efficiency of all oscillators. We can see from the inset in Fig. 3 that transport efficiency of Oscillator-2 is enhanced only because the energy dissipated through Oscillator-1 becomes smaller. This clearly shows that the effect of noise is to create new pathways in the system through which energy can efficiently flow towards an specific site, generally referred to as sink or reaction center.

## Discussion

The results presented here show that noise-assisted transport, an ubiquitous concept that may help us understand efficient energy transport in diverse classical and quantum

systems, can be observed in simple electronic circuit networks. This opens fascinating routes towards new methods for enhancing the efficiency of different energy transport systems, from small-scale RF and microwave electronic circuits to long-distance high-voltage electrical lines. In this way, a specific feature initially conceived in a quantum scenario (environment-assisted quantum transport) has shown to apply as well in classical systems, widening thus the scope of possible quantum-inspired technological applications.

## Methods

In our experiment, three identical  $RLC$  electrical oscillators interact by means of three ideal capacitors. The input energy is injected into a single oscillator (Oscillator-1), and we follow the dynamics of the system by registering independently the voltage of each oscillator (Fig. 1). Noise is introduced in the system by means of random fluctuations in one of the capacitive couplings, where the magnitude of the changes in the capacitance is defined by the noise voltage provided by an arbitrary function generator (Diligent Analog Discovery 410-244P). Relative changes in the capacitance range from 0% (no voltage) to 100% (1 V) in a controllable way. To avoid instabilities, the 1-KHz-frequency noise signal was varied from 0 to 850 mV, which corresponds to a range of the fluctuating capacitance going from  $C_{12}$  to  $(1 \pm 0.85)C_{12}$ .

The voltage of each oscillator contains the information about the stored energy in the system as well as the energy dissipated by each oscillator. A Tektronix MSO4034 oscilloscope (impedance  $1M\Omega$ ) is used to measure these voltages. The voltage signals were extracted from the oscilloscope using a PC-OSCILLOSCOPE interface, which transfers the information through a USB port. Because we are working with stochastic events, measurements were repeated up to 500 times for each noise voltage, and averaged using a MatLab script. The initial input signal is a single pulse, with a pulse duration of 200  $\mu$ s, injected in Oscillator-1 using an Arbitrary Waveform Generator from Agilent 33220A, with a 5 Hz frequency. The high level voltage amplitude is 5 V, while the low-level voltage amplitude is 0 V.

The synthesized transformation of the electronic circuit to an analog computer is described in the Supplementary Information. Basically, the analog computer makes analog simulations of differential equations using electronic components. The input and output

voltages of an electronic circuit correspond to mathematical variables. These voltage variables are therefore representations of the physical variables used in the mathematical model.

An important issue that we would like to point out is that in this work both, the initial conditions and noise, are physically introduced via a voltage signal. This is particularly relevant because in the first case this voltage represents the initial current and, in the second case, the statistical distribution of the noise voltage can be defined independently from the circuit, thus it is not necessary to produce fluctuations in the physical properties of the electrical components—namely resistors, capacitors or inductors—which generally represents a major challenge.<sup>1</sup>

Using the building blocks described in the Supplementary Information, it is possible to synthesize current variables and inductor elements with only resistors, capacitors and voltages. In our experiment, the components employed for the implementation of the corresponding building blocks include metal resistors (1% tolerance), polyester capacitors and general-purpose operational amplifiers (MC1458). A DC power source (BK Precision 1761) was used to generate a  $\pm 12$  V bias voltage for the operational amplifiers. The electronic components of each oscillator were mounted and soldered on a drilled phenolic board (7.5×4.5 cm) to avoid poor contacts.

## References

- <sup>1</sup> Plawsky, J. L. *Transport phenomena fundamentals* (CRC Press, U.S., 2014)
- <sup>2</sup> Mohseni, M., Rebentrost, P., Lloyd, S. & Aspuru-Guzik, A. Environment-assisted quantum walks in photosynthetic energy transfer. *J. Chem. Phys.* **129**, 174106 (2008).
- <sup>3</sup> Plenio, M. & Huelga, S. Dephasing-assisted transport: quantum networks and biomolecules. *New J. Phys.* **10**, 113019 (2008).
- <sup>4</sup> Caruso, F., Chin A. W., Datta A., Huelga S. F. & Plenio, M. B. Highly efficient energy excitation transfer in light-harvesting complexes: The fundamental role of noise-assisted transport. *J. Chem. Phys.* **131**, 105106 (2009).
- <sup>5</sup> Ball, P. Physics of life: The dawn of quantum biology. *Nature* **474**, 272 (2011).
- <sup>6</sup> Lambert, N. *et al.* Quantum biology. *Nat. Phys.* **9**, 10 (2013).
- <sup>7</sup> Huelga, S. F. & Plenio, M. B. Vibrations, Quanta and Biology. *Contemp. Phys.* **54**, 181 (2013).
- <sup>8</sup> Rebentrost, P., Mohseni, M., Kassal, I., Lloyd, S. & Aspuru-Guzik, A. Environment-assisted quantum transport. *New J. Phys.* **11**, 033003 (2009).
- <sup>9</sup> Chin, A., Datta, A., Caruso, F., Huelga, S. & Plenio, M. Noise-assisted energy transfer in quantum networks and light-harvesting complexes. *New J. Phys.* **12**, 065002 (2010).
- <sup>10</sup> Caruso, F., Spagnolo, N., Vitelli, C., Sciarrino, F. & Plenio, M. B. Simulation of noise-assisted transport via optical cavity networks *Phys. Rev. A* **83**, 013811 (2011).
- <sup>11</sup> Mostame, S. *et al.* Quantum simulator of an open quantum system using superconducting qubits: exciton transport in photosynthetic complexes. *New J. Phys.* **14**, 105013 (2012).
- <sup>12</sup> Shabani, A., Mohseni, M., Rabitz, H. & Lloyd, S. Efficient estimation of energy transfer efficiency in light-harvesting complexes. *Phys. Rev. E* **86**, 011915 (2012).
- <sup>13</sup> Kassal, I. & Aspuru-Guzik, A. Environment-assisted quantum transport in ordered systems. *New J. Phys.* **14**, 053041 (2012).

- <sup>14</sup> León-Montiel, R. de J., Kassal, I. & Torres, J. P. Importance of excitation and trapping conditions in photosynthetic environment-assisted energy transport. *J. Phys. Chem. B* **118**, 10588 (2014).
- <sup>15</sup> Kramer, B. & MacKinnon, A. Localization: Theory and experiment. *Rep. Prog. Phys.* **56**, 1469 (1993).
- <sup>16</sup> Hänggi, P. & Marchesoni, F. Artificial Brownian motors: Controlling transport on the nanoscale. *Rev. Mod. Phys.* **81**, 387 (2009).
- <sup>17</sup> León-Montiel, R. de J. & Torres, J. P. Highly efficient noise-assisted energy transport in classical oscillator systems. *Phys. Rev. Lett.* **110**, 218101 (2013).
- <sup>18</sup> Spiechowicz, J., Hänggi, P. & Luczka, J. Brownian motors in the microscale domain: Enhancement of efficiency by noise. *Phys. Rev. E* **90**, 032104 (2014).
- <sup>19</sup> Viciani, S., Lima, M., Bellini, M. & Caruso, F. Observation of Noise-assisted transport in an all-optical cavity-based network. *Phys. Rev. Lett.* **115**, 083601 (2015).
- <sup>20</sup> Biggerstaff, D. N. *et al.* Enhancing quantum transport in a photonic network using controllable decoherence. *Arxiv:1504.06152v1* (2015).
- <sup>21</sup> Caruso, F., Crespi, A., Ciriolo, A. G., Sciarrino, F. & Osellame, R. Fast escape from quantum mazes in integrated photonics. *Arxiv:1501.06438v1* (2015).
- <sup>22</sup> Broome, M. A. *et al.* Discrete single-photon quantum walks with tunable decoherence. *Phys Rev. Lett.* **104**, 153602 (2010).
- <sup>23</sup> Schreiber, A. *et al.* Decoherence and disorder in quantum walks: from ballistic spread to localization. *Phys. Rev. Lett.* **106**, 180403 (2011).
- <sup>24</sup> Svozilík, J., León-Montiel, R. de J. & Torres, J. P. Implementation of a spatial two-dimensional quantum random walk with tunable decoherence. *Phys. Rev. A* **86**, 052327 (2012).
- <sup>25</sup> Rudavskii, Y., Ponedilok, G. & Petriv, Y. Diagonal and off-diagonal types of disorder in condensed matter theory. *Condensed Matter Physics* **3**, 807 (2000).
- <sup>26</sup> León-Montiel, R. de J., Svozilík, J. & Torres, J. P. Generation of a tunable environment for electrical oscillator systems. *Phys. Rev. E* **90**, 012108 (2014).



## **Acknowledgments**

RJLM acknowledges postdoctoral financial support from the University of California Institute for Mexico and the United States (UC MEXUS). MAQJ acknowledges CONACyT for a PhD scholarship. RQT thanks DGAPA-UNAM for financial support through grant IN111614. JLDJ thanks Catedras CONACYT-UNAM. JPT acknowledges support from the Severo Ochoa program (Government of Spain), from the ICREA Academia program (ICREA, Generalitat de Catalunya) and from Fundacio Privada Cellex, Barcelona. JLA wishes to thank CONACYT and DGAPA-UNAM for financial support through grants 167244 and IN106115, respectively.

## **Author Contributions**

RJLM, MAQJ, RQT, JLDJ, JPT and JLA conceived, designed and implemented the experimental setup. RJLM and HMMC provided the theoretical analysis. All authors contributed extensively to the planning, discussion and writing up of this work.

## **Competing Financial Interests**

The authors declare no competing financial interests.

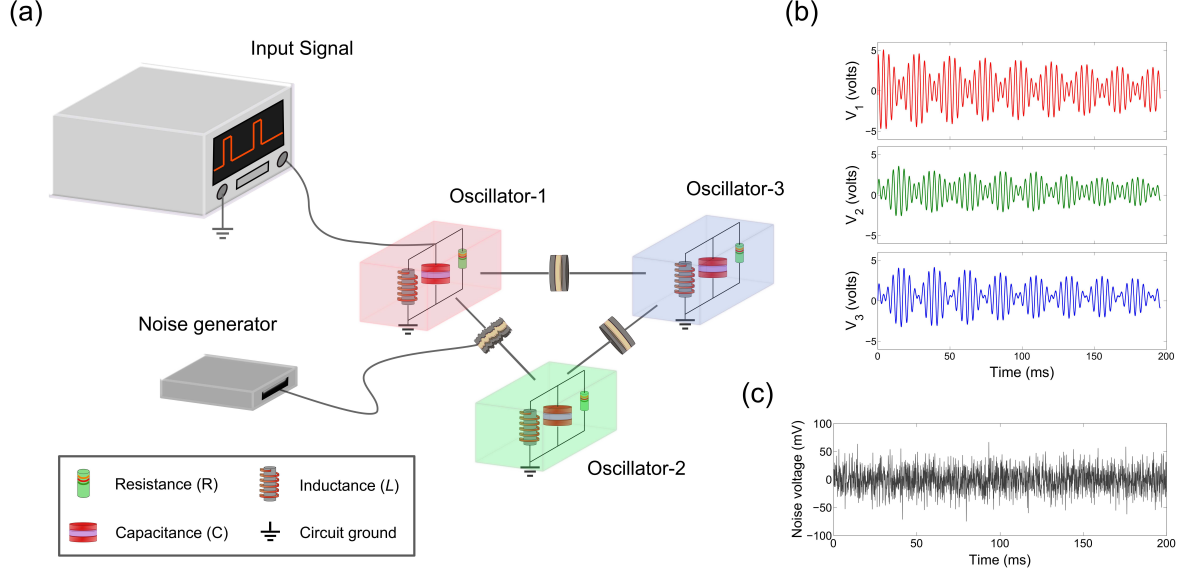


Figure 1: Schematic representation of the experiment and typical dynamic interaction response. (a) Schematic of the network of three capacitively coupled  $RLC$  electrical oscillators. (b) Typical averaged voltage signals measured in each oscillator. (c) Sample of a typical noise signal,  $V_s$ , introduced in the capacitive coupling between the first and second oscillator.

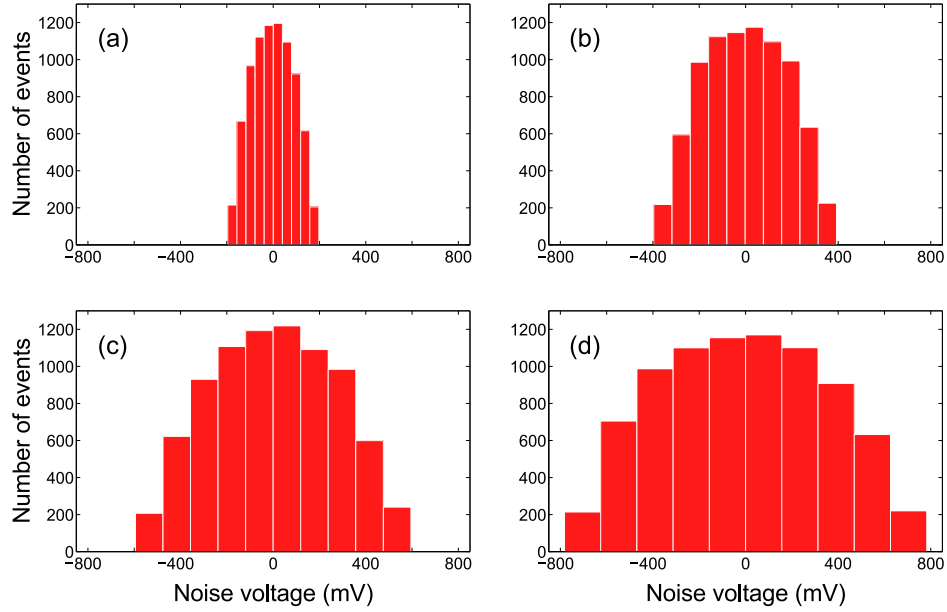


Figure 2: Histograms of noise signals extracted directly from the arbitrary function generator using different values of noise voltage: (a)  $V_s = 200$  mV, (b)  $V_s = 400$  mV, (c)  $V_s = 600$  mV and (d)  $V_s = 800$  mV. Samples were obtained from noise signals with frequency of 1 kHz, within a 1s time-window. Number of events is defined as the number of samples that have the same voltage.

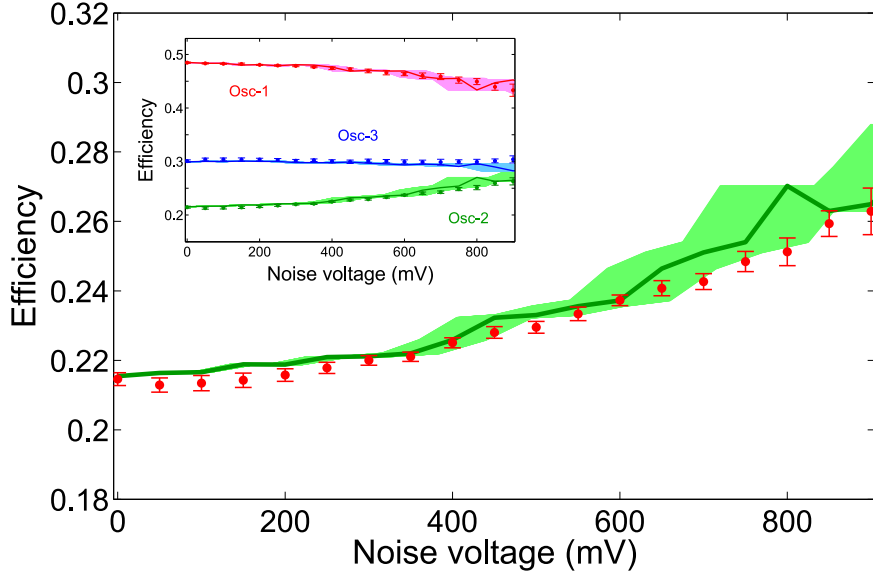


Figure 3: Transport efficiency measured in Oscillator-2 as a function of the noise voltage. Experimental results (dotted line) were obtained by averaging the oscillator's signal over 500 stochastic realizations. The solid line represents the theoretical calculation of transport enhancement using the noise signal extracted directly from the arbitrary function generator. The shaded region represents transport enhancement when noise introduced in the capacitive coupling deviates from the value provided by the function generator by up to 10%. Transport efficiencies of all oscillators are plotted in the inset. Notice that the effect of noise is to rearrange the energy available in the system in order to increase the efficiency of Oscillator-2. Error bars correspond to one standard deviation.

## Supplementary Information: Noise-assisted energy transport in electrical oscillator networks with off-diagonal dynamical disorder

From an experimental point of view, it is challenging to implement a set of interacting  $RLC$  electrical oscillators where certain important defining parameters should change dynamically (noise), even more if the time needed to visualize the dynamics and interactions is large compared with the time-constant imposed by the parasitic attenuation of a real system. In this work we are interested in noise, its implementation, its effects and how to control it. In the first instance, one could attempt to build noisy  $RLC$  oscillators using passive components,<sup>1</sup> but this strategy is in general limited by the large damping coefficient in the inductor, which produces a fast energy extinction, thus hindering the observation of the sought-after effects.

In this supplementary information material, we will show how a network of noisy  $RLC$  oscillators can be electronically implemented by using the working principle of an analog computer.<sup>2,3</sup> Basically, an analog computer makes analog simulations of differential equations using electronic components. The input and output voltage variables of an analog electronic circuit (analog computer) represent physical variables of the mathematical model that is being studied. In this way, in order to simulate any mathematical model in an analog computer, the sequence of mathematical operations involved in the process must be described by means of functional block diagrams. The four functional blocks shown in Fig. 4 are the ones used in our experimental setup to implement the mathematical operations described by Eqs. (1)-(3) of the main manuscript.

Using these functional blocks a complete diagram of the oscillator network can be obtained. Firstly, we build a block diagram for a simple  $RLC$  oscillator [as shown in Fig. 5(a)]. Notice that this diagram corresponds to the synthesis of a typical damped harmonic oscillator equation, with  $\alpha$  and  $\omega$  representing its damping coefficient and frequency, respectively. Once the block diagram is built, each mathematical operation can be electronically implemented with passive linear electrical components and basic inverting configurations of operational amplifiers (OPAMPs).

Figure 5(b) shows the synthesized analog circuit for a typical  $RLC$  electrical oscillator. Notice that the inherent sign inversion is a result of the negative voltage gain of the OPAMP, which must be taken into account in the design. Here,  $R_j$ ,  $C_j$  and  $U_j$  stand for resistors and capacitors, and general-purpose operational amplifiers, respectively. The values of  $R_{f1}$ ,  $R_1$ ,  $R_{i1}$ ,  $C_{i1}$ ,  $R_{i2}$  and  $C_{i2}$  are defined by the oscillator parameters, that is, the resistance ( $R$ ), the inductance ( $L$ ) and the capacitance ( $C$ ). They must satisfy:

$$\frac{R_{f1}}{R_1} = \frac{1}{R}, \quad \frac{1}{R_{i1}C_{i1}} = \frac{1}{C}, \quad \frac{1}{R_{i2}C_{i2}} = \frac{1}{L}, \quad (5)$$

and

$$\frac{R_{f2}}{R_2} = \frac{R_{f2}}{R_3} = \frac{R_{f2}}{R_4} = 1. \quad (6)$$

This last expression [Eq. (2)] implies that the adder ( $U_4$ ) has a unitary amplification factor.

For the experimental setup, we calculated the resistors and capacitors involved in Fig. 5(b) considering the following parameters:  $L = 1$  mH,  $C = 300$   $\mu$ F and  $R = 1$  k $\Omega$ . The combination of these parameters results in an oscillation frequency of  $f = 290.57$  Hz, with a damping coefficient  $\alpha = 3.333$  Hz. To design this configuration we thus take  $R_{i1} = 3$  k $\Omega$ ,  $C_{i1} = 100$  nF,  $R_{i2} = 10$  K $\Omega$ ,  $C_{i2} = 100$  nF,  $R_1 = 1$  k $\Omega$  and  $R_{f1} = 3.3$  k $\Omega$ . By selecting  $R_{f2} = R_2 = R_3 = R_4 = 10$  k $\Omega$ , the condition given by Eq. (2) is satisfied.

To understand how random fluctuations in the coupling between Oscillator-1 and Oscillator-2 are introduced, we first substitute Eq. (3) into Eqs. (1) and (2) of the main manuscript, so we obtain a set of equations given by

$$CV'_1 = -i_1 - \frac{V_1}{R} - C_x(V'_1 - V'_2)(1 + \phi) - C_xV'_1 + C_xV'_3, \quad (7)$$

$$Li'_1 = V_1, \quad (8)$$

$$CV'_2 = -i_2 - \frac{V_2}{R} + C_x(V'_1 - V'_2)(1 + \phi) - C_xV'_2 + C_xV'_3, \quad (9)$$

$$Li'_2 = V_2, \quad (10)$$

$$CV'_3 = -i_3 - \frac{V_3}{R} - 2C_xV'_3 + C_xV'_1 + C_xV'_2, \quad (11)$$

$$Li'_3 = V_3, \quad (12)$$

where we have assumed that the coupling capacitors have the same value, that is,  $C_x = C_{12} = C_{13} = C_{23}$ .

The synthesis of Eqs. (3)-(8) is divided into three parts. In the first part, the *RLC* oscillators are implemented using the electronic circuit in Fig. 5. Each oscillator can be easily implemented using the same electronic circuit that was obtained for the simple *RLC* parallel oscillator. In the second part, the couplings without noise are built using the same methodology, that is, by designing the block diagram and then replacing each block with electronic configurations based on operational amplifiers. Finally, the third part is the design of the noisy coupling. Note that  $CV'_1$  and  $CV'_2$  have the same noisy coupling term thus it is enough to build only one of them. Figures 6(a)-(b) show the functional block diagram of the noisy coupling, as well as its equivalent analog electronic circuit.

To conclude, it is important to point out that, in our system, noise is physically introduced via a voltage signal. This is particularly relevant because its statistical distribution will be defined independently from the circuit, which implies that fluctuations in the physical properties of electronic components—such as resistors, capacitors or inductors—are not required.

## REFERENCES

- <sup>1</sup> León-Montiel, R. de J., Svozilík, J. & Torres, J. P. Generation of a tunable environment for electrical oscillator systems. *Phys. Rev. E* **90**, 012108 (2014).
- <sup>2</sup> The Education and Training Department. Carlson, A., Hannauer, G., Carey, T. & Holsberg, J. (Eds) *Handbook of analog computation* (Electronic Associates Inc., New Jersey, 1967) Associates.
- <sup>3</sup> Johnson, C. L. *Analog computer techniques* (McGraw-Hill, New York, 1963)

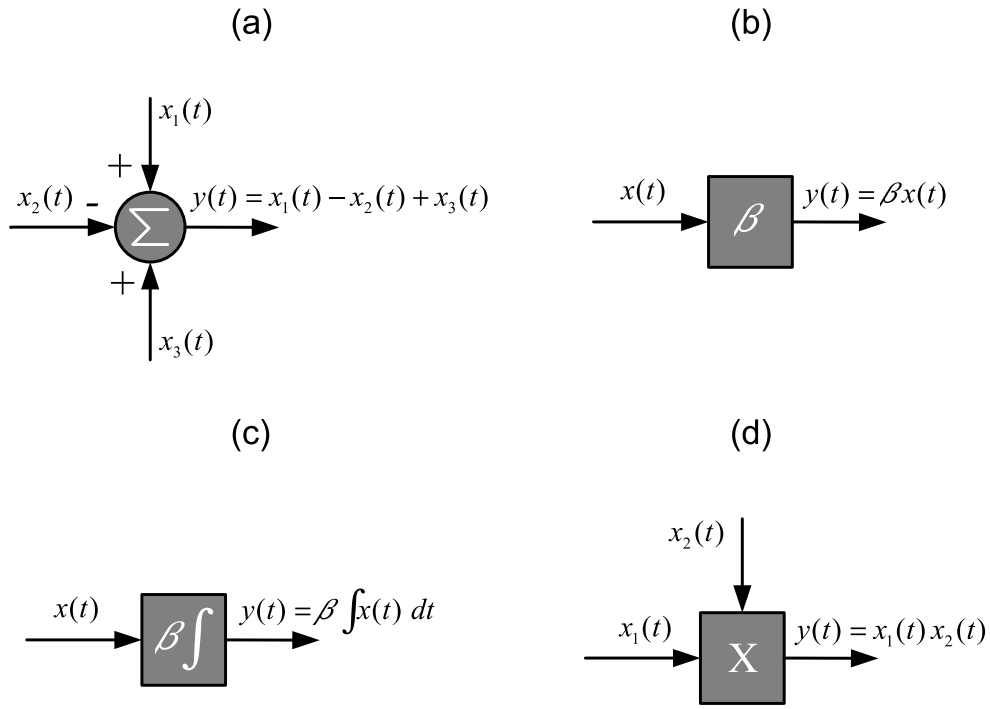


Figure 4: Functional Blocks: (a) adder, (b) multiplication by a constant, (c) integrator and (d) multiplication of signals.

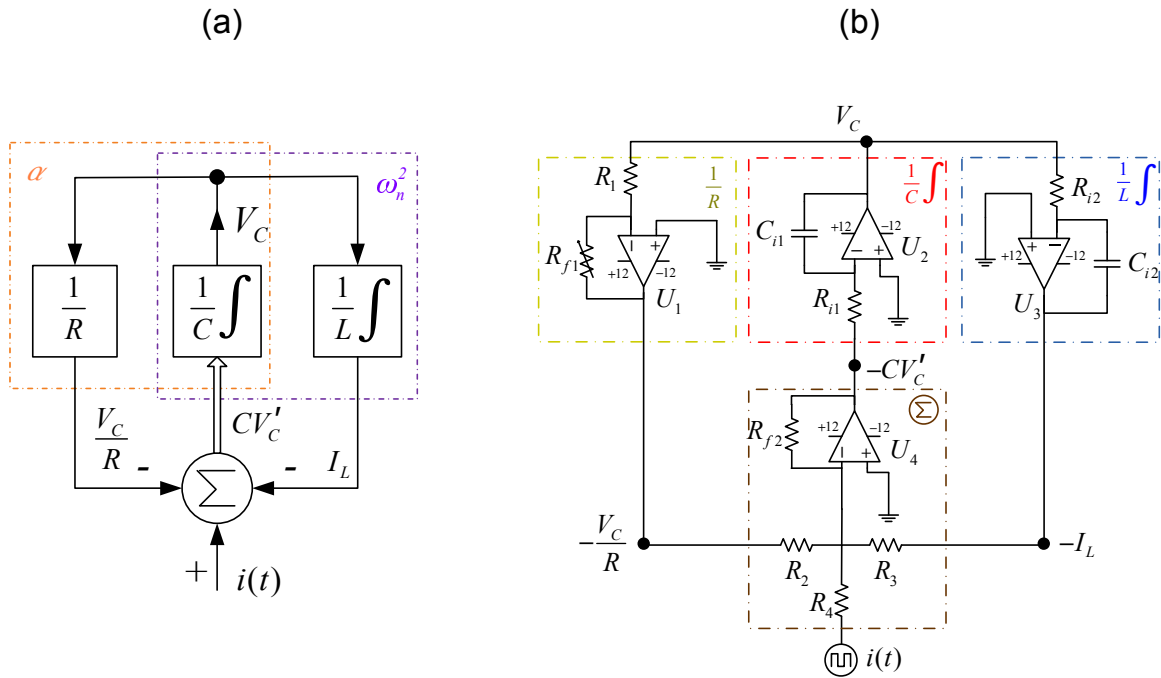


Figure 5: Simple  $RLC$  parallel oscillator (a) Block diagram and (b) Analog electronic circuit.

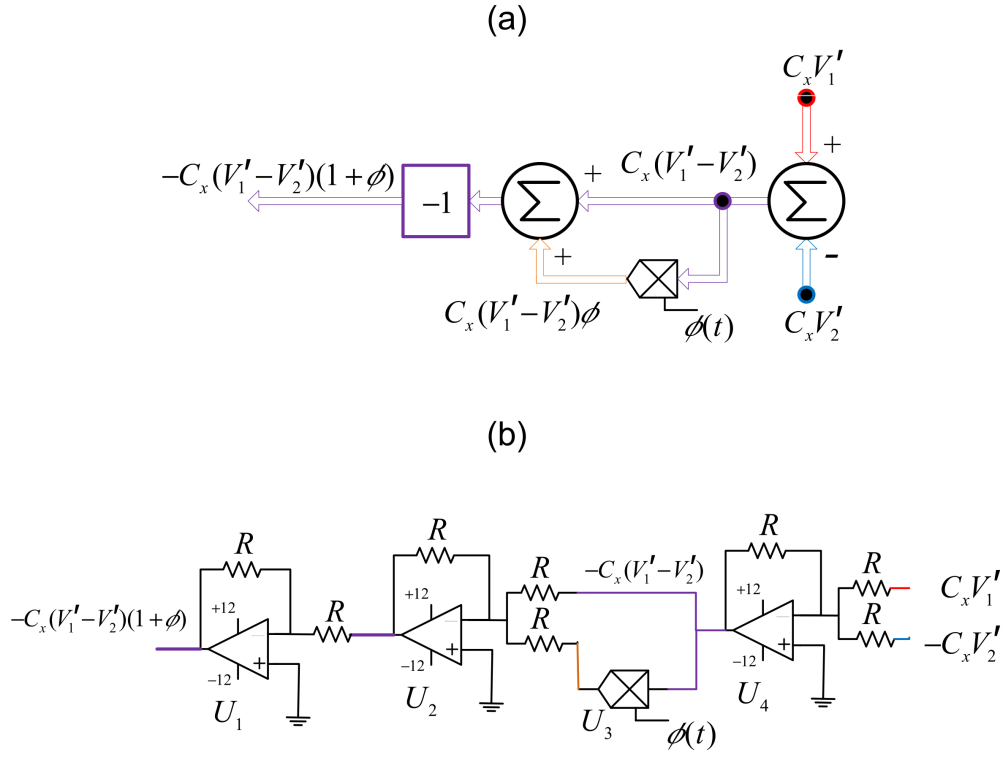


Figure 6: Noisy coupling between Oscillator-1 and Oscillator-2: (a) Block diagram and (b) Analog electronic circuit.